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Finite Element Implementation of a Peridynamic Pitting Corrosion Damage Model

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Abstract: Despite the significant improvements in the understanding of pitting corrosion, many aspects of this phenomenon remain unclear and corrosion rate prediction based on experimental data remains difficult. Experimental measurements of corrosion rates under different electrochemical conditions can be complex and time consuming, and the conclusions are limited to the timescale and the conditions in which experiments have been carried out. In order to overcome these limitations, numerical approaches can be a valuable complement. Hence, in this study a new numerical model based on peridynamics to predict pitting corrosion damage is developed. The developed model is implemented in a commercial finite element software and it allows for the reproduction of realistic pitting morphologies, modelling of microstructural effects such as the presence of intermetallic particles and the reduction of the computational cost of the simulations. In conclusion, the results of this study shows that the peridynamic models can be helpful in failure analysis and design of new corrosion-resistant materials.

Keywords: Pitting; Corrosion; Peridynamics; Damage; Finite Element Method

1. Introduction

Corrosion is a natural process that can be described as the chemical or electrochemical reaction of an engineering material with the environment, which eventually leads to the deterioration of its mechanical properties. There are various types of corrosion mechanisms including stress corrosion cracking (SCC) and pitting corrosion.

Pitting is a localised form of corrosion that leads to the formation of corrosion cavities or pits due to the breakage of the material's passive film. Pitting typically occurs in materials such as stainless steel, aluminium, titanium, copper, magnesium and nickel alloys. Pitting corrosion is the most common, dangerous and destructive type of corrosion in marine and offshore structures (Bhandari et al., 2015).

Despite the significant improvements in the understanding of pitting corrosion, many aspects of this phenomenon remain unclear and corrosion rate prediction based on experimental data remains difficult (Onishi et al., 2012). Experimental measurements of corrosion rates under different electrochemical conditions can be complex and time consuming, and the conclusions are limited to the timescale and the conditions in which experiments have been carried out (Van der Weeën et al., 2014). Moreover, distinguishing the effect of each individual variable becomes extremely difficult (Xiao and Chaudhuri, 2011). In order to overcome these limitations, numerical approaches can be a valuable complement.

Several numerical approaches are available in the literature for the modelling of corrosion damage, and many of them are based on finite element method (Sharland et al., 1989; Walton, 1990; Laycock and White, 2001; Scheiner and Hellmich, 2009). However, the modelling of pit propagation by using finite element method (FEM) is not an easy task since special procedures of moving mesh and remeshing are often needed (Duddu, 2014). Moreover, if remeshing techniques are used, the finite element matrices have to be recalculated for every time step (Vagbharathi and Gopalakrishnan, 2014). In order to overcome these difficulties, an alternative

numerical approach to the FEM can be the use of extended Finite Element Method (XFEM). Vagbharathi and Gopalakrishnan (2014) and Duddu (2014) describe numerical models of corrosion pit propagation in stainless steel based on the XFEM framework and Fick's law of diffusion. An extension to the latter work was done by Duddu et al. (2015), where the Fick's law of diffusion was replaced by the Nernst-Planck equation and crevice corrosion in Al-Mg alloy microstructures was investigated.

Although all the aforementioned models have been quite effective in predicting the corrosion rates reported in the experimental literature, none of these models have been able to predict corrosion subsurface damage, which is a phenomenon documented in the experimental literature (Song et al., 2014). With the aim of filling this gap, a new continuum mechanics formulation, peridynamics (PD), has been recently used to create a model of pitting corrosion (Chen and Bobaru, 2015) based on a modified Nernst-Planck equation. This work was extended in (Chen et al., 2015), where the influence of the pit cover on the pit morphology was investigated.

In this study, a new numerical model based on PD theory for the investigation of pitting corrosion is introduced. This work is based on the PD model of pitting corrosion developed by Chen and Bobaru (2015). However, the numerical approach used in the present study is different and offers the advantage of a reduced computational time by performing the numerical implementation in a commercial finite element software and utilizing implicit time integration. Moreover, the capabilities of the numerical framework have been extended to the modelling of realistic pit morphologies and determining the influence of microstructural features of the material, e.g. intermetallic particles. The ultimate goal of this study is to produce a non-conventional and more effective numerical framework that can be helpful in failure analysis and design of new corrosion-resistant materials.

2. Peridynamics

The governing equations of classical continuum mechanics (CCM) are based on partial differential equations (PDEs) and its mathematical formulation breaks down in the presence of discontinuities such as cracks. This limitation is partially overcome with the adoption of external crack growth criteria based on fracture mechanics. However, this approach presents its own limitations. In light of the limiting assumptions and difficulties of the current approaches, a new mathematical formulation of continuum mechanics was developed by Silling (2000), which is called “peridynamics”.

The governing equations of PD are integro-differential. The PD equations of motion of a generic material point \mathbf{x} can be written as (Silling, 2000)

$$\rho(\mathbf{x})\ddot{\mathbf{u}}(\mathbf{x},t) = \int_{H_{\mathbf{x}}} \mathbf{f}(\mathbf{u}(\mathbf{x}',t) - \mathbf{u}(\mathbf{x},t), \mathbf{x}' - \mathbf{x}) dV_{\mathbf{x}'} + \mathbf{b}(\mathbf{x},t) \quad (1)$$

where $\rho(\mathbf{x})$ and $\ddot{\mathbf{u}}(\mathbf{x},t)$ denote the density and acceleration of the material point \mathbf{x} at time t , respectively. In Eq. (1), $\mathbf{f}(\mathbf{u}(\mathbf{x}',t) - \mathbf{u}(\mathbf{x},t), \mathbf{x}' - \mathbf{x})$ represents the PD force between material points \mathbf{x} and \mathbf{x}' (also called mechanical response function) and the term $\mathbf{b}(\mathbf{x},t)$ is the body force acting on material point \mathbf{x} at time t . $dV_{\mathbf{x}'}$ is the volume associated with material point \mathbf{x}' . According to this new formulation, a material point can interact with other material points not only within its nearest neighbourhood, but also with material points in a larger neighbourhood (Fig. 1).

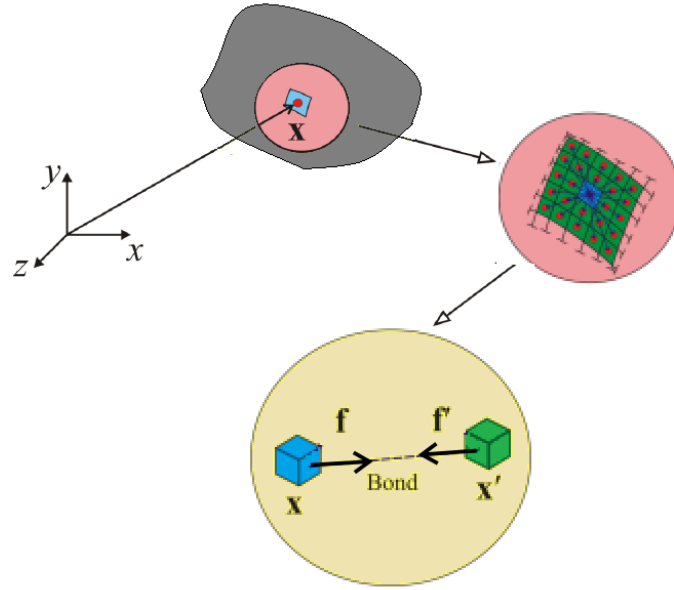


Fig. 1 Peridynamic forces between material points \mathbf{x} and \mathbf{x}' .

It can be assumed that the interaction between material points decreases as the distance between them increases. Therefore, an influence domain, named horizon, $H_{\mathbf{x}}$, can be defined for each material point as shown in Fig. 2.

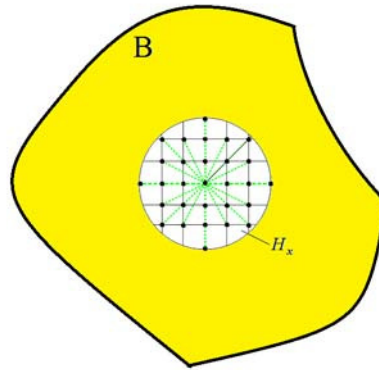


Fig. 2 Peridynamic bonds between material points inside the peridynamic horizon, $H_{\mathbf{x}}$.

The material point \mathbf{x} can only interact with material points within this domain, which are called the “family” of \mathbf{x} . This interaction is called “bond” and its length is simply the distance between the two material points. In the case of an elastic

material, the peridynamic force between material points \mathbf{x} and \mathbf{x}' , can be expressed as

$$\mathbf{f} = c s \frac{\mathbf{y}' - \mathbf{y}}{|\mathbf{y}' - \mathbf{y}|} \quad (2)$$

where \mathbf{y} represents the location of the material point \mathbf{x} in the deformed configuration as shown in Fig. 3, i.e. $\mathbf{y} = \mathbf{x} + \mathbf{u}$, while c is the bond constant which can be related to material constants of CCM as described in Madenci and Oterkus (2014).

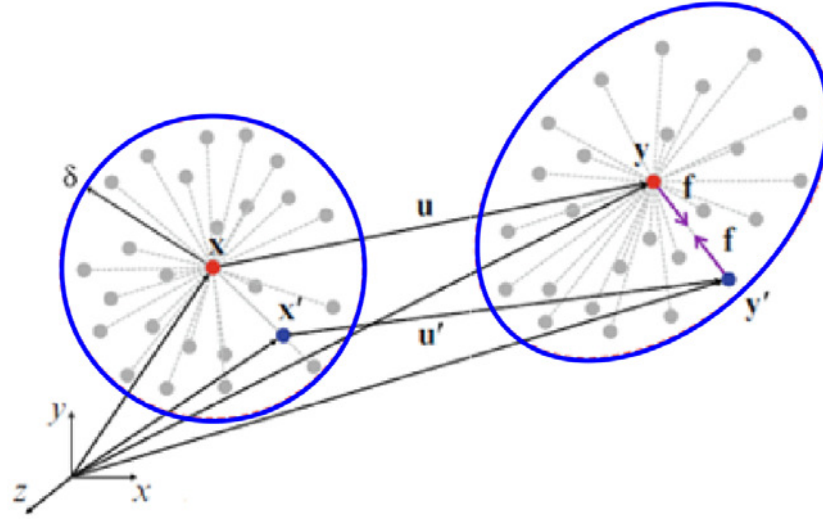


Fig. 3 Peridynamic horizon in the undeformed configuration (left) and the deformed configuration (right)

In Eq. (2), the stretch parameter, s , is defined as

$$s = \frac{|\mathbf{y}' - \mathbf{y}| - |\mathbf{x}' - \mathbf{x}|}{|\mathbf{x}' - \mathbf{x}|} \quad (3)$$

In the case of a brittle material behaviour, the peridynamic force and the stretch relationship are shown in Fig. 4.

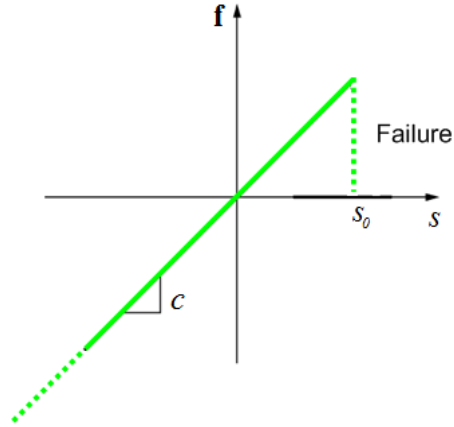


Fig. 4 PD bond behaviour for brittle materials

The parameter s_0 , in Fig. 4, is called critical stretch and if the stretch of a peridynamic bond exceeds this critical value, the peridynamic interaction (bond) is broken. As a result, the peridynamic force between the two material points reduces to zero and the load is redistributed among the other bonds, leading to unguided material failure. Peridynamics can be applicable for different material systems including metals and composites (Oterkus and Madenci, 2012). Moreover, PD framework can be extended to other fields such as thermal (Gerstle et al., 2008), moisture (Oterkus et al., 2014), etc., so that it can be used as a single platform for multiphysics analysis of materials (Amani et al., 2016; De Meo et al., 2016; Oterkus et al., 2013; Oterkus and Madenci, 2014).

3. Peridynamic modelling of pitting corrosion

In this section, PD framework for modelling of pitting corrosion based on the formulation developed by Chen and Bobaru (2015) is presented. The mathematical formulation is based on a modified version of the Nernst-Planck equation (NPE). The NPE is a mass conservation equation that describes the motion of charged chemical species in a fluid under the effect of concentration gradients (diffusion), electric field (migration) and fluid velocity (convection). In this regard, the flux \mathbf{N}_i in $[\text{mol}/(\text{m}^2/\text{s})]$ of the generic species ' i ' can be written as (Gavrilov et al., 2007; Chen and Bobaru, 2015)

$$\mathbf{N}_i = \underbrace{-D_i \nabla C_i}_{\text{diffusion}} - \underbrace{\frac{F}{RT} n_i D_i C_i \nabla \phi}_{\text{electro-migration}} + \underbrace{C_i \mathbf{U}}_{\text{convection}} \quad (4)$$

where D_i in $[\text{m}^2/\text{s}]$ is the diffusion coefficient, C_i is the concentration in $[\text{mol}/\text{m}^3]$ of the species “ i ”, F in $[\text{C}/\text{mol}]$ is the Faraday constant, R in $[\text{J}/(\text{mol K})]$ is the universal gas constant, n_i is the valence number, ϕ in $[\text{V}]$ or $[\text{J}/\text{C}]$ is the electric potential and \mathbf{U} in $[\text{m}/\text{s}]$ is the flow velocity. The conservation of mass can be written as (Gavrilov et al., 2007)

$$\frac{\partial C_i}{\partial t} = -\nabla \cdot \mathbf{N}_i \quad (5)$$

where t in $[\text{s}]$ is time and $\nabla \cdot$ is the divergence operator. The majority of pitting corrosion models available in the literature focuses on the motion of chemical species inside the electrolyte solution. However, recent experimental studies have revealed the existence of a ‘wet region’ at the solid/liquid interface of the corrosion pit, where the motion of chemical species can occur. In other words, the electrolyte solution penetrates the subsurface region of the pit filling the pores of the material, allowing metal dissolution to occur not only at the solid/liquid interface, but also in a subsurface region of the pit.

In order to capture this process, Eq. (5) is used to predict the motion of metal cations Me^{n+} in both the electrolyte solution and the solid with the following simplifications and assumptions:

Motion within the solution

- The electromigration and convective terms in Eq. (4) are included inside the diffusion coefficient, which is now called as effective diffusion coefficient in the liquid D_{ld} .

Therefore, the governing equation used to predict the motion of metal cations within the electrolyte solution is given by the following modified version of the Nernst-Planck equation as

$$\frac{\partial C_{MI}}{\partial t} = -\nabla \cdot (-D_{ld} \nabla C_{MI}) = D_{ld} \nabla^2 C_{MI} \quad (6)$$

where, C_{MI} refers to the metal concentration inside the electrolyte solution, which is considered to be 1M NaCl in this study. The effective diffusion coefficient in the liquid can be obtained from the study reported in Kuo and Landolt (1975), where rotating disk electrode experiments of iron dissolution in chloride solutions are described and the electromigration effect is taken into consideration empirically as

$$D_{ld} = 8.5 \cdot 10^{-10} \text{ m}^2 / \text{s} \quad (7)$$

Motion within the solid

- The velocity of the fluid diffusing inside the pores of the metal is neglected and the convection term in Eq. (4) is not considered.
- The effect of the electromigration term in Eq. (4) is included inside the diffusion coefficient, which is now called as effective diffusion coefficient in the solid D_{sd} .

The governing equation used to describe the motion of metal cations within the solid is given by the following modified version of the Nernst-Planck equation as

$$\frac{\partial C_{MI}}{\partial t} = -\nabla \cdot (-D_{sd} \nabla C_{MI}) = D_{sd} \nabla^2 C_{MI} \quad (8)$$

where, in this case, C_{MI} refers to the metal concentration inside the solid.

As shown in Eq. (8), the flux of metal cations is linearly proportional to the effective diffusivity. According to Faraday's laws of electrolysis, the amount of

metal dissolution N_a in [mol/s] is linearly proportional to the anodic current density, which can be expressed as

$$i_a = nFN_a = nF \frac{dM_{ol}}{dt} \quad (9)$$

where M_{ol} expressed in [mol] is the number of moles of the material. Therefore, in order to obey Faraday's laws of electrolysis, the effective diffusion coefficient in the solid is expressed as a function of the overpotential η as

$$D_{sd} = D_{sd}(\eta) = D_{sd}(0)e^{\frac{\alpha n F \eta}{RT}} \quad (10)$$

In Eq. (10), the term $D_{sd}(0)$ represents the value of the diffusion coefficient in the solid when the overpotential is null, which can be found through calibration against experimental polarisation data. Eq. (10) can also be written as

$$\eta = \beta_a \cdot \log \left(\frac{D_{sd}(\eta)}{D_{sd}(0)} \right) \quad (11)$$

Moreover, the anodic current density in Eq. (9) can also be expressed as a function of the moving speed of the solid/liquid interface, i.e. the pit base, v_{int} [m/s] as

$$i_a = nFN_a = nFC_{solid}v_{int} \quad (12)$$

where $F = 9.64859 \cdot 10^4$ C/mol and C_{solid} in [mol/m³] is the concentration of metal cations in the solid. As described in Scheiner and Hellmich (2009), in the case of stainless steel grade 304 (304 SS), i.e. the material considered in this study, C_{solid} can be calculated by considering the mass density of 304 SS and the molar masses of its main components Fe, Ni and Cr as

$$C_{solid} = 143,000 \text{ mol / m}^3 \quad (13)$$

A similar approach can be used for the calculation of the valence number of 304 SS:

$$n = 2.19 \quad (14)$$

As mentioned earlier, the corrosive solution considered in this study is 1M NaCl at temperature $T = 298.15$ K and the related saturation concentration value of dissolved stainless steel metal cations, C_{sat} , can be inferred from X-ray analysis (Isaacs et al., 1995) as

$$C_{sat} = 5,100 \text{ mol / m}^3 \quad (15)$$

As described in Chen and Bobaru (2015), the calibration procedure for the calculation of $D_{sd}(0)$ consists of the following steps:

1. A small initial value is chosen for the diffusion coefficient in the solid $D_{sd}(\eta)$ to make sure that corrosion will be in activation-controlled regime.
2. The simulation is launched and the velocity v_{int} of the solid/liquid interface, i.e. metal dissolution front, is recorded.
3. Eq. (12) is used to calculate the current density, i_a .
4. Once i_a is known, it is possible to calculate the corresponding overpotential value by using the experimental polarization diagram reported in Chen and Bobaru (2015).
5. Eq. (11) can now be used to calculate $D_{sd}(0)$ considering the following value of the Tafel constant: $\beta_a = 45.1$ mV (Chen and Bobaru, 2015).

$D_{sd}(0) = 7.352 \cdot 10^{-17} \text{ m}^2/\text{s}$ is the outcome of this calibration procedure, which is in agreement with Chen and Bobaru (2015).

When the concentration of metal ions in the liquid reaches the saturation value C_{sat} , a salt film precipitates at the liquid/solid interface. The concentration of metal

ions in the liquid cannot be greater than the saturation value. When the concentration value of the generic node is greater than C_{sat} , then the node is considered to be in solid phase. On the contrary, when the concentration value is smaller than C_{sat} , then the node is considered to be in liquid phase.

The peridynamic governing equation for metal dissolution can be written as

$$\dot{C}_{MI}(\mathbf{x}, t) = \int_{H_{\mathbf{x}}} f_d(C_{MI}(\mathbf{x}, t), C_{MI}(\mathbf{x}', t), \mathbf{x}', \mathbf{x}, t) dV_{\mathbf{x}'} \quad (16)$$

where $\dot{C}_{MI}(\mathbf{x}, t)$ is the time derivative of metal ions concentration associated with the generic material point \mathbf{x} . In Eq. (16), the peridynamic function $f_d(C_{MI}(\mathbf{x}, t), C_{MI}(\mathbf{x}', t), \mathbf{x}', \mathbf{x}, t)$ is called metal dissolution response function and it is defined as

$$f_d = d_{MI} \frac{C_{MI}(\mathbf{x}', t) - C_{MI}(\mathbf{x}, t)}{|\mathbf{x}' - \mathbf{x}|} \quad (17)$$

in which the peridynamic metal ions diffusion bond constant d_{MI} can be expressed in terms of the effective diffusion coefficient as (Chen and Bobaru, 2015)

$$(2D) \quad d_{MI} = \frac{4 \cdot D_{eff}}{\pi \cdot \delta^2 \cdot h} \quad (18)$$

where h in [m] refers to the thickness of the body and D_{eff} in [m²/s] is the effective diffusion coefficient which can be calculated as suggested in Chen and Bobaru (2015) as

$$D_{eff} = \begin{cases} D_{sd} & \text{if } C_i > C_{sat} \text{ and } C_j > C_{sat} \\ D_{ld} & \text{if } C_i < C_{sat} \text{ and } C_j < C_{sat} \\ \frac{2 \cdot D_{sd} \cdot D_{ld}}{D_{sd} + D_{ld}} & \text{if } C_i \leq C_{sat} \text{ and } C_j \geq C_{sat} \\ \frac{2 \cdot D_{sd} \cdot D_{ld}}{D_{sd} + D_{ld}} & \text{if } C_i \geq C_{sat} \text{ and } C_j \leq C_{sat} \end{cases} \quad (19)$$

The PD model of pitting corrosion consists of mechanical bonds overlapped by diffusion bonds. The former bonds are aimed to capture the subsurface mechanical damage reported in recent experimental studies (Song et al. 2014). For this purpose, a damage index, $d(\mathbf{x}, t)$, is defined for each node as

$$d(\mathbf{x}, t) = \frac{N_f}{N_{tot}} \quad (20)$$

where N_f and N_{tot} are the number of failed bonds and the total number of bonds attached to the node at \mathbf{x} , respectively. At the beginning of the simulation, all the nodes belonging to the solid body have a concentration value $C_{MI}(\mathbf{x}, t)$ equal to C_{solid} , which means no corrosion. Therefore, in this condition, all the mechanical bonds connected to \mathbf{x} are intact and its damage $d(\mathbf{x}, t)$ index is null. On the contrary, when $C_{MI}(\mathbf{x}, t)$ becomes smaller than C_{liquid} , the node at \mathbf{x} changes its phase from solid to liquid, which means complete metal dissolution. Therefore, in this condition, all the mechanical bonds connected to \mathbf{x} are broken and the damage index has a unit value. As suggested in Chen and Bobaru (2015), the damage index $d(\mathbf{x}, t)$ can therefore be written as a function of the nodal concentration $C_{MI}(\mathbf{x}, t)$ as

$$d(\mathbf{x}, t) = \frac{C_{solid} - C_{MI}(\mathbf{x}, t)}{C_{solid} - C_{sat}} \quad (21)$$

Each bond has a certain probability P_r to be broken in the current time step. In each time step, P_r can be calculated as (Chen and Bobaru, 2015)

$$P_r = \frac{1}{1-d_{t-1}} \cdot \left(\frac{\Delta C_i}{C_{solid} - C_{sat}} \right) \quad (22)$$

where d_{t-1} represents the value of the damage index at the previous time step, and ΔC_i is the difference in nodal concentration between the previous time step and the current time step. Once P_r is calculated, a random number in the range $[0,1]$ is generated for each bond. If the random number is smaller than P_r , then the mechanical bond is broken and the value of the damage index is updated.

4. Numerical solution method

As mentioned earlier, the PD equation of motion is integro-differential and, in general, it cannot be solved analytically. Therefore, numerical techniques for time and spatial integration are usually needed to solve the PD governing equations.

Concerning spatial integration, a meshless scheme and the collocation method can be used. The domain can be divided in smaller parts (Fig. 5), where each part has an associated volume and integration material point, which is also called as collocation point. In the case of a uniform grid, the distance among the material points is same in all directions.

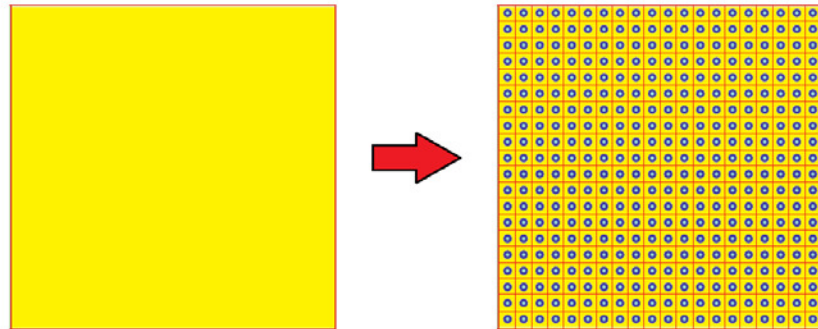


Fig. 5 Discretisation of PD domain in 2-Dimensional configurations.

For a generic material point \mathbf{x} , the spatial integration is performed only over the part of the body that is contained within the horizon of the particle \mathbf{x} . Therefore, for a generic PD particle \mathbf{x}_i , the discretised form of Eq. (1) can be written as

$$\dot{C}_{MI}(\mathbf{x}, t) = \sum_j^M f_d \left(C_{MI}(\mathbf{x}_j, t) - C_{MI}(\mathbf{x}_i, t), \mathbf{x}_j - \mathbf{x}_i \right) V_j \quad (23)$$

where M is the number of family members of particle \mathbf{x}_i . Concerning the time integration, implicit time integration scheme has been utilized in this study. Eq. (23) can be solved by using either explicit or implicit time integration schemes. The advantage of an implicit time integration scheme is the numerical stability of the method regardless of the value of the time step size. This is particularly important when modelling corrosion, since the timescale of the problem can range from seconds to years. In this study, a novel approach, based on the analogy between heat transfer and metal ions diffusion, is used to exploit the implicit solver implemented in the commercial finite element analysis software ANSYS 14.5.

Thus, the PD model of pitting corrosion can be represented by the following thermal finite elements available in ANSYS:

- MASS 71: Thermal mass elements are used to represent the PD particles. The only degree of freedom is the temperature, which, in our model, corresponds to the metal ions concentration [mol/m³].
- COMBIN 14: Thermal spring elements are utilized to represent the PD bonds. These elements are used to connect the thermal mass elements to each other. As for the MASS 71 elements, the only degree of freedom is the temperature.

The expression for the real constant d_{ANSYS} of the generic COMBIN 14 element can be written as

$$(2D) \quad d_{ANSYS} = \frac{d_{MI} \cdot s_{cf} \cdot V \cdot v_c \cdot V}{l^2} \quad (24)$$

where l in [m] represents the length of the bond and V in [m³] represents the volume a material point. In Eq. (24), s_{cf} and v_c are surface and volume corrections, respectively, which are explained in Madenci and Oterkus (2014).

5. Numerical results

In this section, three different cases are presented. In the first case, pitting corrosion evolution in a stainless steel plate is simulated for validation purposes. In the second case, more realistic pit morphologies are generated. Finally, the effect of intermetallic particles on pitting corrosion evolution is investigated.

5.1 Pitting corrosion in a stainless steel plate

In order to validate the current approach, a 2D plate of dimensions 0.1 mm x 0.1 mm and thickness of 1 μ m is considered. As shown in Fig. 6, the top edge of the plate is damaged and the initial pit is filled with 1M NaCl aqueous solution, in which the concentration of metal ions is initially null. The volume occupied by the solution is assumed to be much larger than that of the metal, i.e. austenitic stainless steel grade 304. This assumption allows to represent the bulk of the corrosive solution with a zero-concentration boundary condition throughout the simulation.

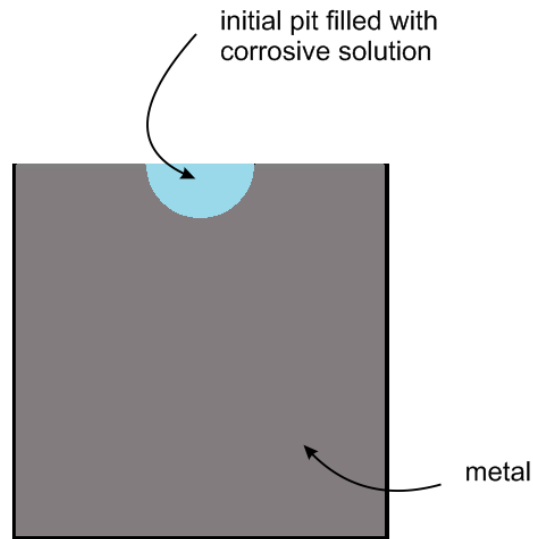


Fig. 6 2D pitting corrosion model: initial pit (blue colour) and metal (grey colour)

The material data can be found in Section 3. The radius of the initial pit is $10\text{ }\mu\text{m}$. The plate is discretised with 100×100 PD nodes and the resulting value of grid spacing and horizon's radius are $\Delta = 1\text{ }\mu\text{m}$ and $\delta = 4\text{ }\mu\text{m}$, respectively. At the beginning of the simulation, all the nodes belonging to the metal have a concentration of metal ions equal to C_{solid} . The appropriate time step size is found by performing a convergence analysis. As shown in Fig. 7, a time step size of 0.5 s can be considered as a suitable value. The results obtained by using the present model are in qualitative agreement with the numerical results presented in Chen and Bobaru (2015), where a time step size of 0.05 s is used (see Fig. 8).

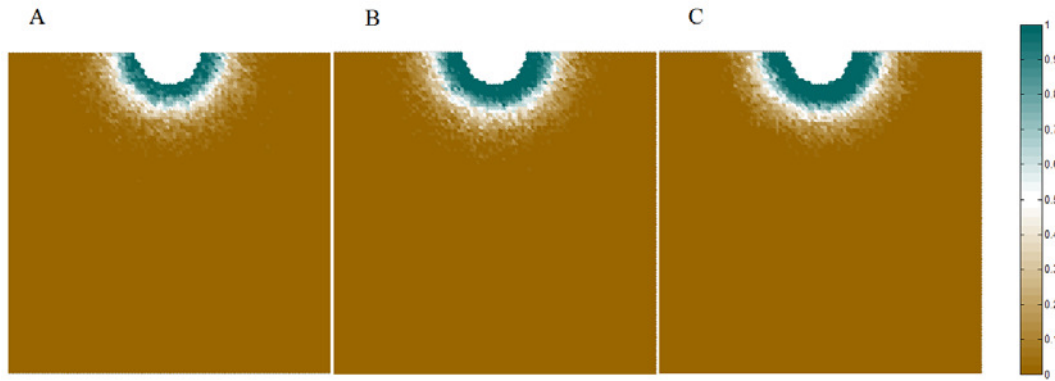


Fig. 7 Time step size convergence analysis for the damage index map at time = 20 s and $\eta = 0.2\text{ V}$: $dt = 5\text{ s}$ (A), $dt = 0.5\text{ s}$ (B), $dt = 0.05\text{ s}$ (C)

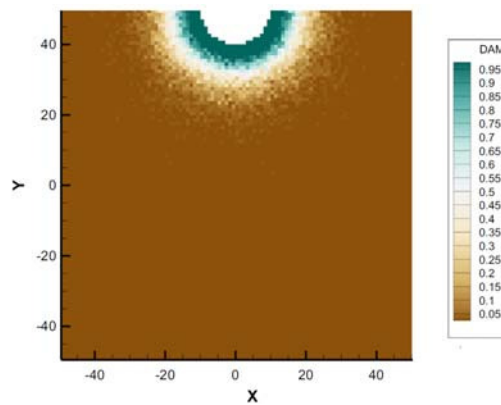


Fig. 8 Damage index map at time = 20 s for the case with $\eta = 0.2$ V reported in Chen and Bobaru (2015)

In Figs. 7 and 8, only the particles in solid phase are shown. The liquid/solid interface does not advance within the first 20 s, but a layer constituted by solid particles with damage index value equals to one is formed at the interface. As explained earlier, if a generic particle has a damage index value equals to one, it means that all of its bonds have been broken. Thus, in this case, the particle is totally disconnected from the rest of the body and the solution can penetrate underneath. In order to mimic this phenomenon, the approach used in Chen and Bobaru (2015) is adopted. According to this approach, the phase of the node is changed from solid to liquid when $d = 1$ even if $C > C_{sat}$. As shown in Fig. 9, the results produced by the present model are in quantitative agreement with the study reported in Chen and Bobaru (2015).

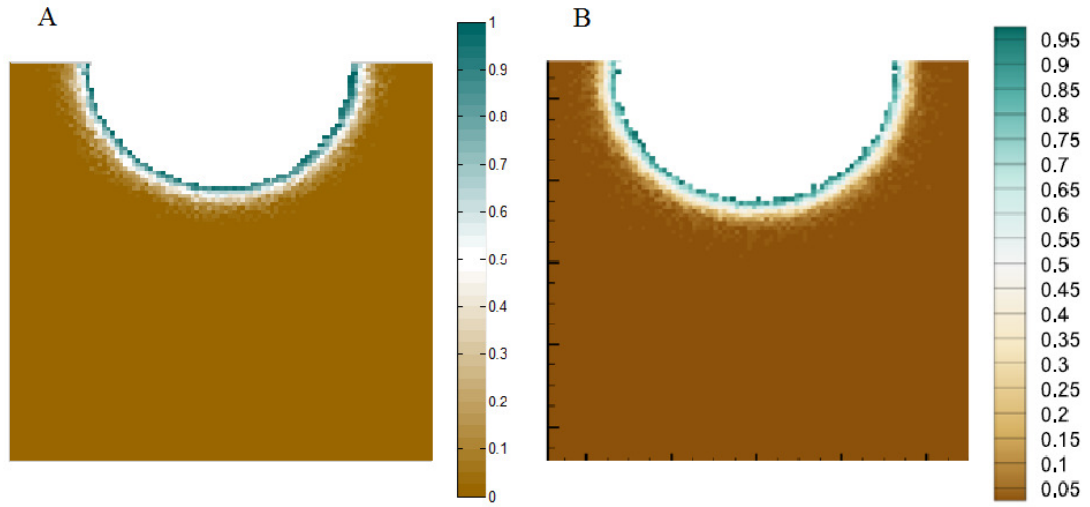


Fig. 9 Damage index map at time = 20 s for the case with $\eta = 0.2$ V: present numerical model (A) and other numerical model (Chen and Bobaru, 2015) (B)

5.2 Realistic pit morphologies

In reality, corrosion pits can have complex geometries, which are often very different from the simple case considered in the previous section. The PD model of pitting corrosion can be extended to model realistic pit morphologies. There may be

several different approaches that can be used in this regard. The procedure adopted in this study, which is used during the preparation of the model, consists of the following two steps:

1. Selection of the region of the domain where the pit is expected to propagate.
2. For all the nodes that do not belong to this region, an artificial effective coefficient of diffusion in the metal D_{sd}^* is used. The ratio between D_{sd}^* and D_{sd} is always smaller than 1 and is called “coefficient of pit morphology” c_{pm} :

$$c_{pm} = \frac{D_{sd}^*}{D_{sd}} < 1 \quad (25)$$

The following boundary condition is applied to the blue region of length $6 \mu\text{m}$ in Fig. 10 as $C(x, t) = 0$.

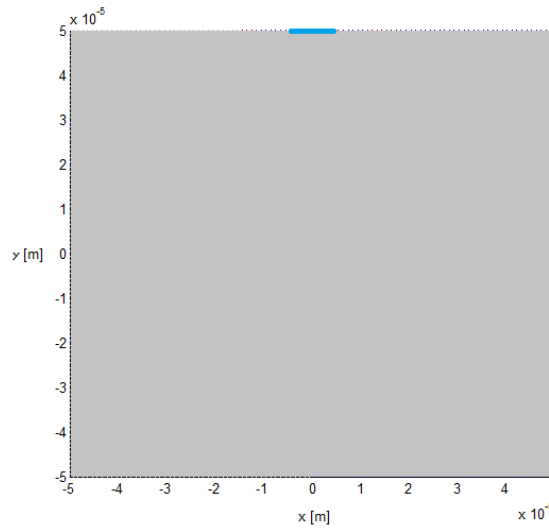


Fig. 10 PD model for the prediction of realistic pit morphologies: metal (grey colour) and initial liquid/solid interface (blue colour)

As shown in Fig. 11, if the morphology of the pit is known a priori, then the PD model is able to reproduce the expected shape of the pit cavity. The material

properties and numerical parameters used for the simulations in Fig. 11 are the same as those used in previous section. The only difference is the introduction of c_{pm} , which is assumed to be 1%.

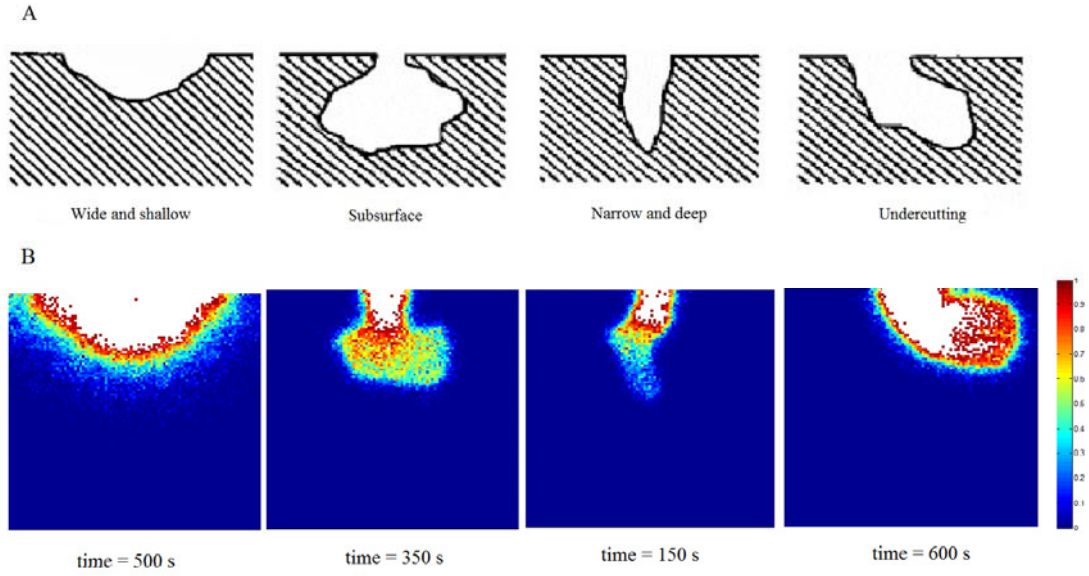


Fig. 11 Realistic pit morphologies (Hoeppepner, 2011) (A), damage index map obtained with the present numerical model (B)

5.3 Pitting corrosion in a cluster of cathodic intermetallic particles

It is well-known that the metallurgical characteristics of the material can affect the evolution of corrosion pits. In this analysis, pitting propagation in a cluster of cathodic intermetallic particles (IMPs) is investigated. The material properties and numerical parameters are the same as those used in the previous section. IMPs are created through node selection during the generation of the model and they are represented by the application of the following boundary condition to all those nodes that belong to the IMPs domain as $C(x, t) = C_{solid}$.

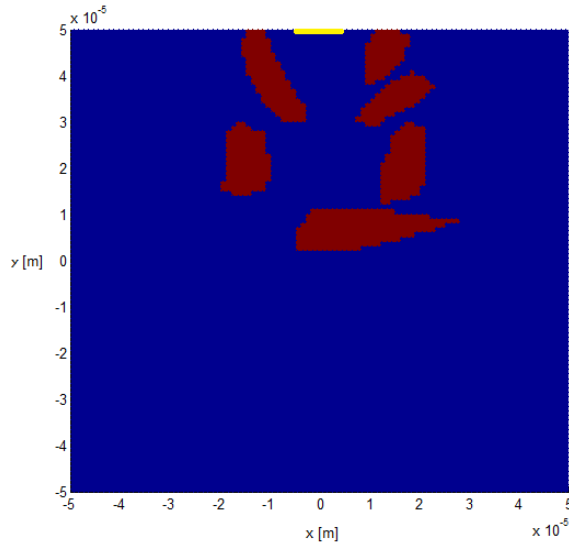


Fig. 12 PD model of pitting corrosion in a cluster of IMPs: metal (blue colour), corrosive solution (yellow colour) and IMPs (red colour)

As for the analyses described in previous sections, the following boundary condition is applied to the yellow region of length $6 \mu\text{m}$ in Fig. 12 as $C(x,t) = 0$. As shown in Fig. 13, pit damage is channelled by the cluster of IMPs and the lower part of the plate is partially shielded by the IMP at the bottom of the cluster. In order to better visualise the pit cavity, only solid nodes are shown in Fig. 13.

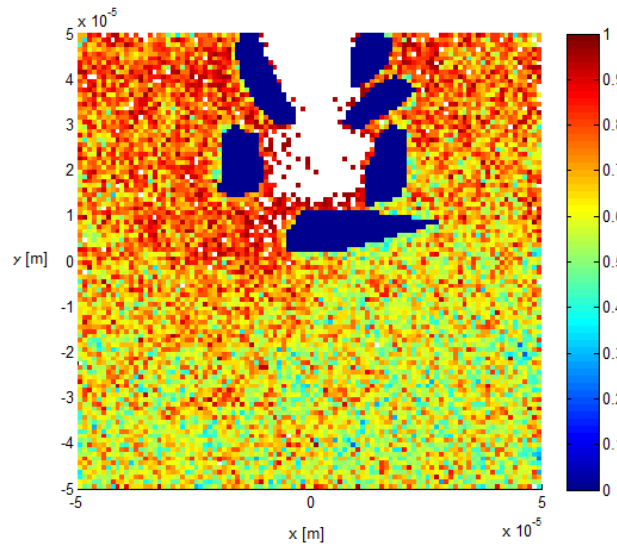


Fig. 13 Pit evolution in a cluster of IMPs. Damage index map at time = 2500 s: no dissolution (blue colour) and close to dissolution (red colour)

6. Conclusions

The ultimate goal of this study was to use peridynamics to produce a non-conventional and more effective numerical framework that can be helpful in failure analysis and design of new corrosion-resistant materials. For this purpose, a new numerical model based on peridynamics to predict pitting corrosion damage was developed by extending an existing model to allow for 1) the reproduction of realistic pitting morphologies, 2) modelling of microstructural effects such as the presence of intermetallic particles and 3) the reduction of the computational cost of the simulations. Three different cases were considered to validate the developed model and demonstrate its capabilities. In the first case, pitting corrosion evolution in a stainless steel plate was simulated and a suitable time step size was determined after performing a convergence analysis. In the second case, realistic pit morphologies were obtained by using the current framework. Finally, the effect of intermetallic particles on pitting corrosion evolution was investigated and it was found that the intermetallic particles have a significant effect on pitting corrosion damage evolution.

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